# ANALYSIS OF CYCLONE PRESSURE DROP <br> Lingjuan Wang, Graduate Research Assistant <br> Calvin B. Parnell, Jr Ph.D., P. E., Regent Professor Bryan W. Shaw, Ph.D., Assistant Professor <br> Department of Agricultural Engineering <br> Texas A\&M University <br> College Station, TX 


#### Abstract

A new method to analysis cyclone pressure drop is reported. The frictional pressure loss is the primary pressure loss in a cyclone. The air stream travel distance is a function of cyclone diameter. The frictional pressure loss is independent of a cyclone diameter, therefore cyclone total pressure loss is independent of cyclone diameter.


## Introduction

Cyclone separators use centrifugal force to separate suspended particulate matter from the air stream. A cyclone consists of a cylindrical upper body with a conical lower section and a smaller center cylinder that extends from the top to just below the entrance through which the relatively clean air is discharged (Figure 1.). The dust-laden air stream enters tangentially at the top of the barrel portion of the cyclone and travels downward into the cone portion of the cyclone forming an outer vortex. The increasing air velocity in the outer vortex results in an increased centrifugal force on the particles separating them from the air stream. When the air reaches the bottom of the cone, an inner vortex is created reversing the air stream direction in the center of the cone. The air stream leaves the cyclone through the center cylinder at the top of the cyclone while the particulate falls into the dust collection hopper attached to the bottom of the cyclone.

The operation of a cyclone is relatively simple but not completely understood. Extensive work has been completed to determine how the cyclone dimensions and its operating conditions affect its performance. However, the engineering data associated with air and particle flow patterns in the cyclone are difficult to accurately measure. The current cyclone design theory is based partly on theoretical analysis and partly on empirical models. The goal of this work is to develop a sound science description of the operation of a cyclone that can be used to facilitate engineering design with a minimum of empirical data.

The three parameters used to evaluate the cyclone performance are emission concentrations, collection efficiency as a function of particle size and energy consumption (pressure drop) as a function of inlet velocity. Cyclone pressure drop associated with the operation of the cyclone is a major factor to be considered in the design of a cyclone collection system. Many models have been developed to determine drop such as Shepherd and Lapple equation (1939), Stairmand equation (1949, 1951), First equation (1950) and Stern equation (1977). However, the equations are either empirical models or involve variables and dimensionless parameters not easily accounted for in practical applications. It is known that cyclone pressure drop is dependent on the cyclone design and its operating parameters such as inlet velocity. The empirical models cannot be used for all the cyclone designs as new cyclone technology and new cyclone designs are developed. Further theoretical research is needed to scientifically evaluate the cyclone performance including predicting cyclone pressure drop.

## Cyclone Theory

The air flow pattern in a cyclone rotational field is rather complex. It can be characterized by three velocity components such as tangential velocity, axial velocity and radial velocity. Ter Linden (1949) first measured the

National Cotton Council, Memphis TN
details of the flow field in a cyclone and described the air flow tangential velocity distribution and the pressure distribution in a cyclone as given by the Figure 2.

After the air stream enters the cyclone, it will spiral downward because of centrifugal and frictional forces. In the outer vortex, the total air flow velocity consists of tangential velocity, radial velocity and axial velocity. The tangential velocity is the dominant velocity component. It also determines the centrifugal force applied to the air stream. The tangential velocity can be described as follows:

$$
\begin{equation*}
V_{t} * r^{n}=C \tag{1}
\end{equation*}
$$

where:
$\mathrm{V}_{\mathrm{t}}=$ tangential velocity,
$\mathrm{r}=$ rotational radius,
$\mathrm{n}=$ flow pattern factor,
$\mathrm{n}=0.5 \sim 0.8$ (in outer vortex)
$\mathrm{n}=0$ at the boundary of inner vortex and outer vortex
$\mathrm{n}=-1$ in inner vortex
$\mathrm{C}=$ constant. (guangda $\mathrm{Ma}, 1983$ )
The tangential velocity increases with the decreasing of the rotational radius in the outer vortex. It increases to the maximum at the boundary of the outer vortex and inner vortex. In the inner vortex the tangential velocity decreases as the rotational radius decrease. In the inner vortex, the relationship of the tangential velocity and the rotational radius can be modeled by the following equation:

$$
\begin{equation*}
\frac{V_{t}}{r}=\omega=C \tag{2}
\end{equation*}
$$

where:
$\mathrm{V}_{\mathrm{t}}=$ tangential velocity,
$\omega=$ angular velocity,
$\mathrm{C}=$ constant .

The pressure distribution in a cyclone varies with different cyclone designs. The rotation of the air-flow establishes a pressure field. The highest pressure is located at the wall of the cyclone. The pressure decreases with the reduction of the rotational radius because of the increasing tangential velocity. In the radial direction, there is a significant pressure drop caused by the change of the tangential velocity, since the increasing of the tangential velocity causes a large radial acceleration. This pressure change in the radial direction can be expressed by the following equation:

$$
\begin{equation*}
\frac{d P}{d r}=\rho * \frac{V_{t}^{2}}{r} \tag{3}
\end{equation*}
$$

where:
$d P=$ change of the pressure
$d r=$ change of the rotation radius
$\mathrm{V}_{\mathrm{t}}=$ tangential velocity
$\mathrm{R}=$ rotational radius, and
$\rho=$ air density

The solution of the equation 3 is:

$$
\begin{equation*}
\mathrm{P}=\rho^{*} \mathrm{~V}_{\mathrm{t}}^{*} \ln (\mathrm{r})+\mathrm{C} \tag{4}
\end{equation*}
$$

This solution gives the pressure distribution in the cyclone flow field.

## Analysis of the Cyclone Pressure Drop

The cyclone pressure drop is a function of the cyclone dimensions and its operating conditions. Shepherd and Lapple $(1939,1940)$ determined the optimum dimensions of cyclones based on the body diameter. Such general relationships have allowed for a simplified analysis. The definition of the cyclone dimensions is shown in the Figure 3.

The pressure drop over a cyclone is caused by the area changes, the wall friction, change of the flow direction and the dissipation in the vortex finder (outlet tube). Because of these effects, the cyclone pressure drop is mainly composed of the following parts:
(1) $\Delta \mathrm{P}_{1}$--- the entry pressure loss in the tangential inlet duct.
(2) $\Delta \mathrm{P}_{2}$--- the frictional loss along the total spiral travel distance in the outer vortex, caused by the wall friction.
(3) $\Delta \mathrm{P}_{3}$--- the pressure loss caused the change of the air flow direction from the outer vortex to the inner vortex (approximately "U" turn).
(4) $\Delta \mathrm{P}_{4}--$ the entry pressure loss from the point that air enters the inner vortex to the point that air leaves the outlet tube. To determine this part of pressure loss, an assumption is made that the inner vortex forms an imaginary solid cylinder whose outer surface is the boundary between the inner vortex and the outer vortex. This cylinder starts at the point that air enters the inner vortex and ends at the bottom of the outlet tube. The diameter of this inner vortex cylinder is 0.6 times the outlet tube diameter (Ter Linden, 1949, Guangda Ma, 1983).

## Entry Pressure Loss in the Tangential Inlet Duct (DP ${ }_{1}$ )

In this part, the pressure loss is caused by the inlet duct area changes. It can be determined as follows:

$$
\begin{equation*}
\Delta P_{1}=C^{*} V P_{i} \tag{5}
\end{equation*}
$$

where:
$\Delta \mathrm{P}_{1}=$ dynamic pressure loss in the inlet duct,
$\mathrm{C}=$ dynamic loss constant, and
$\mathrm{VP}_{\mathrm{i}}=$ cyclone inlet velocity pressure

## Frictional Pressure Loss ( $\Delta \mathbf{P}_{2}$ )

This part of pressure loss is the pressure loss in the cyclone outer vortex caused by the friction of the gas/surface wall. In the outer vortex, air stream flows in a downward spiral through the cyclone. It may be considered as the air stream travels in a imaginary spiral tube with diameter $D_{s}$ and length $L$ (the air stream travel distance in the outer vortex). It can be determined as follows:

$$
\begin{equation*}
\Delta P_{2}=f * \frac{L}{D_{S}} * V P_{s i} \tag{6}
\end{equation*}
$$

where:
$\Delta \mathrm{P}_{2}=$ the frictional pressure loss,
$\mathrm{f}=$ friction factor, dimensionless,
$\mathrm{L}=$ total travel distance,
$\mathrm{D}_{\mathrm{s}}=$ the imaginary spiral tube diameter, and
$\mathrm{VP}_{\text {si }}=$ air stream inlet velocity pressure in the imaginary spiral tube.

The imaginary spiral tube diameter can be approximated by the following equation:

$$
\begin{equation*}
D_{s}=\left(\frac{D_{c}-D_{e}}{2}\right) \tag{7}
\end{equation*}
$$

where:
$D_{s}=$ the imaginary spiral tube diameter,
$D_{c}=$ cyclone barrel diameter, and
$D_{e}=$ cyclone outlet tube diameter.
The travel distance ( L ) consists of the travel distance in the zone $\mathrm{Z}_{1}\left(\mathrm{~L}_{1}\right)$ plus the travel distance in the zone $\mathrm{Z}_{2}\left(\mathrm{~L}_{2}\right)$.

## 

The air-flow in the zone $Z_{1}$ only has two velocity components: tangential velocity $\left(\mathrm{V}_{\mathrm{t}}\right)$ and the axial velocity $\left(\mathrm{V}_{\mathrm{z}}\right)$ (see Figure 3.). The tangential velocity and axial velocity in this zone are constant. As a result, the travel distance in this zone can be determined by the following equation:

$$
\begin{align*}
& L_{1}=\int_{0}^{t} 1\left|\vec{V}_{t 1}+\vec{V}_{z 1}\right| d t \\
& =\int_{0}^{z} 1 \sqrt{V_{t 1}^{2}+V_{z 1}^{2}} \frac{d z}{V_{z 1}} \tag{8}
\end{align*}
$$

where:
$L_{1}=$ travel distance in the zone $Z_{1}$,
$\mathrm{t}_{1}=$ travel time in the zone $\mathrm{Z}_{1}$,
$Z_{1}=$ the total height of the zone $Z_{1}$,
$\mathrm{V}_{\mathrm{t} 1}=$ tangential velocity in the zone $\mathrm{Z}_{1}\left(\mathrm{~V}_{\mathrm{t} 1}=\mathrm{V}_{\mathrm{i}}\right)$, and
$\mathrm{V}_{\mathrm{z} 1}=$ axial velocity in the zone $\mathrm{Z}_{1}$.
Because the air flow rate $(\mathrm{Q})$ in a cyclone is constant, the axial velocity can be calculated by the following equation:

$$
\begin{align*}
& V_{z 1}=\frac{2 V_{i}}{3 \pi} \\
& =\frac{32 V_{i}}{39 \pi} \quad\left(\text { for } \mathrm{D}_{\mathrm{e}}=\mathrm{D}_{\mathrm{c}} / 2 \text { design }\right)  \tag{9}\\
& \left(\text { for } \mathrm{D}_{\mathrm{e}}=\mathrm{D}_{\mathrm{c}} / 1.6 \text { design }\right)
\end{align*}
$$

The equations (8) and (9) can be used to calculate $L_{1}$ as follows:

$$
\begin{align*}
& L_{1}=\frac{3 * \pi * Z_{1}}{2} * \sqrt{1+\left(\frac{2}{3 \pi}\right)^{2}} \quad\left(\text { for } \mathrm{D}_{\mathrm{e}}=\mathrm{D}_{\mathrm{c}} / 2 \text { design }\right) \\
& =\frac{39 * \pi * Z_{1}}{32} * \sqrt{1+\left(\frac{32}{39 \pi}\right)^{2}} \quad\left(\text { for } \mathrm{D}_{\mathrm{e}}=\mathrm{D}_{\mathrm{c}} / 1.6 \text { design }\right)
\end{align*}
$$

## Travel Distance in the Cone Portion (zone $\mathbf{Z}_{2}$ ) $--\mathbf{L}_{2}$

In the zone $\mathrm{Z}_{2}$, the total velocity consists of three components $\left(\mathrm{V}_{\mathrm{t}}, \mathrm{V}_{\mathrm{z}}, \mathrm{V}_{\mathrm{t}}\right)$ (see Figure 3). The travel distance in this zone is calculated based on those three velocity components.

$$
\begin{align*}
& L_{2}=\left.\int_{0}^{t}\right|_{V} \vec{V}_{t 2}+\vec{V}_{z 2}+\vec{V}_{r 2} \mid d t \\
& \quad=\int_{0}^{z} 2 \sqrt{V_{t}^{2}+V_{z}^{2}+V_{r}^{2}} \frac{d z}{V_{z}} \tag{11}
\end{align*}
$$

where:
$\mathrm{L}_{2}=$ travel distance in the zone $\mathrm{Z}_{2}$,
$\mathrm{t}_{2}=$ travel time in the zone $\mathrm{Z}_{2}$,
$\mathrm{Z}_{2}=$ the total height of the zone $\mathrm{Z}_{2}$,
$\mathrm{V}_{\mathrm{t}}=$ tangential velocity in the zone $\mathrm{Z}_{2}$,
$\mathrm{V}_{\mathrm{z}}=$ axial velocity in the zone $\mathrm{Z}_{2}$, and
$\mathrm{V}_{\mathrm{r}}=$ radial velocity.
In the zone $Z_{2}$, there is an axial acceleration caused by the change of the rotational radius (see Figures 3\&4). The axial velocity at I-I cross section $\left(\mathrm{V}_{21}\right)$ is determined by the equation 9 and $\mathrm{V}_{22}$ can be determined by the equation 12 :

$$
\begin{align*}
V_{z 2} & =\frac{8 * V_{i}}{\pi} \\
= & \left(\text { for } \mathrm{D}_{\mathrm{II}}=\mathrm{D}_{\mathrm{c}} / 4 \text { design }\right), \\
& \frac{2 * V_{i}}{\pi} \quad\left(\text { for } \mathrm{D}_{\mathrm{II}}=\mathrm{D}_{\mathrm{c}} / 2 \operatorname{design}\right)
\end{align*}
$$

The following equation can be used to determine the axial velocity as the Z cross section:

$$
\begin{equation*}
V_{z}=\frac{V_{z 2}-V_{z 1}}{Z_{2}} * Z+V_{z 1} \tag{13}
\end{equation*}
$$

where:
$\mathrm{V}_{\mathrm{z}}=$ the axial velocity at Z cross section,
$\mathrm{V}_{\mathrm{z1}}=$ the axial velocity at I-I cross section,
$\mathrm{V}_{22}=$ the axial velocity at II-II cross section,
$\mathrm{Z}_{2}=$ the total height of the cone portion, and
$Z=$ the high distance from I-I cross section to $Z$ cross section.

The radial velocity $V_{r}$ is determined by the $V_{z}$ (see Figure 5)

$$
\begin{equation*}
\left|\stackrel{\rightharpoonup}{V}_{r}\right|=\operatorname{tg} \theta *\left|\vec{V}_{z}\right| \tag{14}
\end{equation*}
$$

where:
$\mathrm{V}_{\mathrm{r}}=$ the radial velocity at Z cross section,
$\mathrm{V}_{\mathrm{z}}=$ the axial velocity at Z cross section, and $\theta=$ cyclone cone angle.

$$
\operatorname{tg} \theta=\left(\frac{D_{c}-D_{I I}}{2}\right) Z_{2}
$$

$$
\left.=\frac{3 D_{c}}{8 Z_{2}} \quad \text { (for } \mathrm{D}_{\mathrm{II}}=\mathrm{D}_{\mathrm{c}} / 4 \text { design }\right)
$$

The travel distance $L_{2}$ in the zone $Z_{2}$ can be calculated by combining equations (11) through (15).

$$
L_{2}=\frac{3^{*} \pi^{*} Z_{2}}{22}\left[\begin{array}{cc}
s & 1  \tag{16}\\
- & -\frac{a}{2} \\
a & \ln \left|\frac{a+s}{a-s}\right|
\end{array}\right]_{z=0}^{z=z_{2}}
$$

where:
$\mathrm{L}_{2}=$ travel distance in the zone $\mathrm{Z}_{2}$,
$\mathrm{a}=$ mathematical calculation factor for integration,

$$
a=\sqrt{\frac{V_{t}^{2}}{1+t g^{2} \theta}}=\sqrt{\frac{V_{i}^{2}}{1+t g^{2} \theta}}
$$

$\mathrm{s}=$ integrating transformation,

$$
s=\sqrt{a^{2}+V_{z}^{2}}=V_{i} \sqrt{\frac{1}{1+\operatorname{tg}^{2} \theta}+\left(\frac{22 * Z+2 * Z_{2}}{3 * \pi * Z_{2}}\right)^{2}}
$$

The total travel distance L is as:

$$
\begin{gather*}
L=\frac{3}{2} * \pi * Z_{1} * \sqrt{1+\left(\frac{2}{3 \pi}\right)^{2}}+ \\
\frac{3 * \pi * Z_{2}}{22}\left[\begin{array}{cc}
s & 1 \\
- & -\frac{\ln }{a} \quad 2 \\
2 & a+s \\
a-s \mid
\end{array}\right]_{z=0}^{z=z_{2}} \tag{17}
\end{gather*}
$$

The Pressure Loss Caused by the Change of the Flow Direction $\left(\Delta P_{3}\right)$
When the air stream travels to the bottom of the cyclone, the inner vortex reverses the flow direction. The directional change is approximately $180^{\circ}$. The pressure loss caused by this change may be considered as two $90^{\circ}$ elbow losses. However, the air travels to the bottom of the cyclone as a spiral and then reverses the direction in the inner vortex. The directional change is not as sharp as the flow in the 90- degree elbow, therefore, the fitting loss factor should be smaller than that used to calculate the 90-degree elbow pipe.

$$
\begin{equation*}
\Delta \mathrm{P}_{3}=2 * \Delta \mathrm{P}_{90}=2 * \mathrm{~K}_{90} * \mathrm{VP}_{\mathrm{II}} \tag{18}
\end{equation*}
$$

where:
$\Delta \mathrm{P}_{3}=$ pressure loss caused by directional change,
$\mathrm{K}_{90}=$ fitting loss factor, dimensionless, and
$\mathrm{VP}_{\mathrm{II}}=$ axial velocity pressure at the bottom of the cyclone.

## Entry Pressure Loss at the Inner Vortex Entrance ( $\boldsymbol{\Delta} \mathbf{P}_{4}$ )

The assumption has been made that the rotation of the inner vortex forms a imaginary solid cylinder which starts at the bottom of the cyclone and ends at the bottom of the outlet tube. The pressure loss from the entrance of the imaginary cylinder through the outlet tube can be considered as entry loss (see Figure 6).

$$
\begin{equation*}
\Delta \mathrm{P}_{4}=\mathrm{C} * \mathrm{VP}_{\mathrm{II}} \tag{19}
\end{equation*}
$$

where:
$\Delta \mathrm{P}_{4}=$ entry pressure loss at the entrance of the inner vortex,
C = dynamic loss constant, and
$\mathrm{VP}_{\mathrm{II}}=$ axial velocity pressure at the bottom of the cyclone.

## Sample Pressure Drop Calculation and Discussions

## 1D3D @ $\mathbf{V}_{i}=3200 \mathrm{fpm}$

Assume: standard air $\rho=0.075 \mathrm{lb} / \mathrm{ft}^{3}, \mathrm{VP}_{\mathrm{i}}=\left(\mathrm{V}_{\mathrm{i}} / 4005\right)^{2}=0.638$ in $\mathrm{H}_{2} \mathrm{O}$
(1) $\Delta \mathrm{P}_{1}=\mathrm{C}^{*} \mathrm{VP}_{\mathrm{i}}=1^{*} 0.638=0.638$ in $\mathrm{H}_{2} \mathrm{O}$
(2) $\Delta \mathrm{P}_{2}$ : For 1D3D: $\operatorname{tg} \theta=1 / 8, \mathrm{a}=0.99 \mathrm{~V}_{\mathrm{i}}, \mathrm{S}_{z=0}=\mathrm{V}_{\mathrm{i}}, \mathrm{S}_{z=22}=2.73 \mathrm{~V}_{\mathrm{i}}$ $\mathrm{L}=\mathrm{L}_{1}+\mathrm{L}_{2}=4.82 \mathrm{D}_{\mathrm{c}}+5.14 \mathrm{D}_{\mathrm{c}}=9.96 \mathrm{D}_{\mathrm{c}}$
$\mathrm{D}_{\mathrm{s}}=\mathrm{D}_{\mathrm{c}} / 4$
$\mathrm{V}_{\mathrm{si}} *\left(\pi * \mathrm{D}_{\mathrm{s}}{ }^{2} / 4\right)=\mathrm{V}_{\mathrm{i}} *\left(\mathrm{D}_{\mathrm{c}}{ }^{2} / 8\right)$
$\mathrm{V}_{\mathrm{si}}=(8 / \pi) * \mathrm{~V}_{\mathrm{i}}=8153 \mathrm{fpm}$
$\mathrm{VP}_{\mathrm{si}}=(8153 / 4005)^{2}=4.14$ in $\mathrm{H}_{2} \mathrm{O}$
$\Delta \mathrm{P}_{2}=\mathrm{f} *\left(\mathrm{~L} / \mathrm{D}_{\mathrm{s}}\right) * \mathrm{VP}_{\text {si }}=3.3$ in $\mathrm{H}_{2} \mathrm{O}$
$\underline{\Delta \mathrm{P}_{1}}+\Delta \mathrm{P}_{2}=0.638+3.3=3.938$ inch $\mathrm{H}_{2} \mathrm{O}$
The previous research (Askew, 1993) suggested that 1D3D cyclone pressure drop at design velocity was approximately 4.5 inch $\mathrm{H}_{2} \mathrm{O}$. The above sample calculation indicates that cyclone entry pressure loss $\left(\Delta \mathrm{P}_{1}\right)$ and frictional loss in the outer vortex $\left(\Delta \mathrm{P}_{2}\right)$ are the major pressure losses for 1D3D cyclone.

The following observations are made:

- In the 1D3D barrel portion, the travel distance $\left(\mathrm{L}_{\mathrm{t}}\right)$ is $4.82 \mathrm{D}_{\mathrm{c}}$. The number of turns in this portion can be approximated by the following equation:

$$
\begin{equation*}
N=\frac{L_{1}}{\pi * D_{c}}=\frac{4.82 D_{c}}{\pi * D_{c}}=1.53 \tag{20}
\end{equation*}
$$

According to the previous research (Parnell, 1996), there are approximately six turns in the 1D3D. If there are 1.5 turns in the barrel portion, therefore there are approximately 4.5 turns in the cone portion.

- $\frac{L}{D_{S}}=\frac{9.96 D_{c}}{\underline{D_{c}}}=39.84$

4
Equations (6) and (21) demonstrate that the frictional pressure drop is independent of the cyclone diameter. The cyclone entry pressure drop is also independent of the cyclone diameter. Therefore, the total 1D3D cyclone pressure-drop should be constant regardless of the cyclone size.

## 2D2D @ $\mathbf{V}_{i}=3000 \mathrm{fpm}$

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Assume: standard air \(\rho=0.075 \mathrm{lb} / \mathrm{ft}^{3}, \mathrm{VP}_{\mathrm{i}}=\left(\mathrm{V}_{\mathrm{i}} / 4005\right)^{2}=0.561\) in \(\mathrm{H}_{2} \mathrm{O}\)
    (1) \(\Delta \mathrm{P}_{1}=\mathrm{C}^{*} \mathrm{VP}_{\mathrm{i}}=1 * 0.561=0.561\) in \(\mathrm{H}_{2} \mathrm{O}\)
    (2) \(\Delta \mathrm{P}_{2}\) : For 1D3D: \(\operatorname{tg} \theta=3 / 16, \mathrm{a}=0.983 \mathrm{~V}_{\mathrm{i},} \mathrm{S}_{z=0}=\mathrm{V}_{\mathrm{i}}, \mathrm{S}_{z=z 2}\)
        \(=2.73 \mathrm{~V}_{\mathrm{i}}\)
            \(\mathrm{L}=\mathrm{L}_{1}+\mathrm{L}_{2}=9.63 \mathrm{D}_{\mathrm{c}}+3.45 \mathrm{D}_{\mathrm{c}}=13.08 \mathrm{D}_{\mathrm{c}}\)
            \(\mathrm{D}_{\mathrm{s}}=\mathrm{D}_{\mathrm{c}} / 4\)
            \(\mathrm{V}_{\mathrm{si}} *\left(\pi * \mathrm{D}_{\mathrm{s}}{ }^{2} / 4\right)=\mathrm{V}_{\mathrm{i}} *\left(\mathrm{D}_{\mathrm{c}}{ }^{2} / 8\right)\)
            \(\mathrm{V}_{\mathrm{si}}=(8 / \pi) * \mathrm{~V}_{\mathrm{i}}=7639 \mathrm{fpm}\)
            \(\mathrm{VP}_{\mathrm{si}}=(7639 / 4005)^{2}=3.64\) in \(\mathrm{H}_{2} \mathrm{O}\)
        \(\Delta \mathrm{P}_{2}=\mathrm{f} *\left(\mathrm{~L} / \mathrm{D}_{\mathrm{s}}\right) * \mathrm{VP}_{\mathrm{si}}=3.8\) in \(\mathrm{H}_{2} \mathrm{O}\)
\(\underline{\Delta \mathrm{P}_{1}}+\Delta \mathrm{P}_{2}=0.561+3.8=4.371\) inch \(\mathrm{H}_{2} \mathrm{O}\)
```

The cyclone entry pressure loss $\left(\Delta \mathrm{P}_{1}\right)$ and air flow frictional loss in the outer vortex $\left(\Delta \mathrm{P}_{2}\right)$ are the major pressure losses for 2D2D cyclone.

Similar observations are made for the 2D2D cyclone.

- In the 2D2D barrel portion, the travel distance $\left(\mathrm{L}_{1}\right)$ is $9.63 \mathrm{D}_{\mathrm{c}}$. The number of turns in this portion is three. Six turns were observed in the 2D2D. Therefore, there are three turns in the cone portion of a 2D2D cyclone.
- For 2D2D design, $\frac{L}{D_{S}}=\frac{13.08 D_{C}}{\underline{D_{C}}}=52.32$

4

The frictional pressure drop is independent of the cyclone diameter. As was the case with the 1D3D cyclone, the entry pressure drop is also independent of the cyclone diameter. Again, the total 2D2D cyclone pressure-drop should be constant regardless of the cyclone size.

## Problem

There is a problem in the process used to calculate the cyclone frictional pressure drop addressed in this paper. $\left(\mathrm{L} / \mathrm{D}_{\mathrm{s}}\right)$ is used to obtain the major pressure loss in a cyclone. Here, L is the flow travel distance in the outer vortex, and it is also the length of the imaginary spiral tube. $\mathrm{D}_{\mathrm{s}}$ is the imaginary spiral tube diameter. It is treated as a constant. As a matter of fact, $\mathrm{D}_{\mathrm{s}}$ is not constant. It varies with the rotational radius. The smaller the diameter, the higher the pressure drop. Thus, errors were introduced by using a constant stream flow diameter. For this reason the calculated pressure drop for the 1D2D design is larger than that for the 1D3D in the given example. To reduce this error, the stream diameter will varied in future research.

## Summary

Cyclone frictional pressure loss is the major pressure loss in a cyclone. A new assumption has been made to calculate this pressure loss. Air flow in the outer vortex is considered as traveling in a downward spiral tube. The travel distance was calculated based on the three velocity components. The travel distance is the length of the imaginary spiral tube. The diameter of the imaginary spiral tube is a function of the cyclone diameter and the diameter of the cyclone outlet tube. The length of the stream tube divided by its diameter is a constant for a certain cyclone design. The cyclone pressure drop is independent of the cyclone size (diameter).

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Figure 1. Schematic flow diagram of a cyclone.


Figure 2. The tangential velocity and pressure distribution in a cyclone.

$\xrightarrow[\substack{V_{t} \\ \mathbf{V} \text { (tangential) }}]{\substack{V_{r} \\ V_{z}}} \mathbf{Z}$ (axial)
(1) Velocities in the outer vortex Of $Z_{1}$ Portion

$$
\mathrm{v}_{\mathrm{t}} \not \downarrow_{\mathrm{v}_{\mathrm{z}}}
$$

(2) Velocities in the outer vortex
Of $Z_{2}$ Portion


Figure 3. Cyclone dimensions and the velocity distributions.


Figure 4. Axial velocity change zone.


Figure 5. The radial velocity in the $\mathrm{Z}_{2}$ zone.


Figure 6. The imaginary inner vortex cylinder.

